

# CMOS-Compatible Protonic Programmable Resistor Based on Phosphosilicate Glass Electrolyte for Analog Deep Learning

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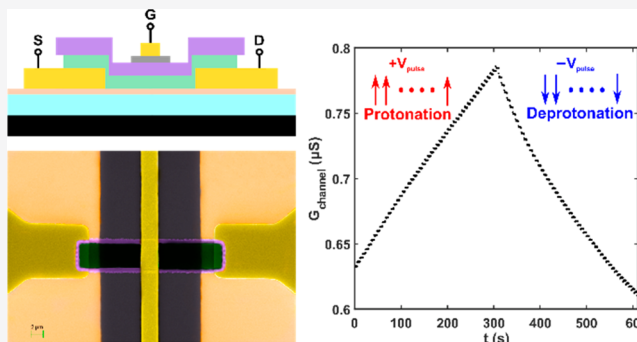
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**ABSTRACT:** Ion intercalation based programmable resistors have emerged as a potential next-generation technology for analog deep-learning applications. Proton, being the smallest ion, is a very promising candidate to enable devices with high modulation speed, low energy consumption, and enhanced endurance. In this work, we report on the first back-end CMOS-compatible nonvolatile protonic programmable resistor enabled by the integration of phosphosilicate glass (PSG) as the proton solid electrolyte layer. PSG is an outstanding solid electrolyte material that displays both excellent protonic conduction and electronic insulation characteristics. Moreover, it is a well-known material within conventional Si fabrication, which enables precise deposition control and scalability. Our scaled all-solid-state three-terminal devices show desirable modulation characteristics in terms of symmetry, retention, endurance, and energy efficiency. Protonic programmable resistors based on phosphosilicate glass, therefore, represent promising candidates to realize nanoscale analog crossbar processors for monolithic CMOS integration.

**KEYWORDS:** analog computing, doped silicon dioxide films, proton intercalation, programmable resistors



## 1. INTRODUCTION

The success of deep learning in classifying and clustering representations of data at multiple levels of abstraction has fundamentally changed how information is processed.<sup>1</sup> However, conventional digital architectures face increasing difficulties in supporting the heavy computational workloads required to train state of the art deep neural networks (DNNs).<sup>2</sup> The pressing need for faster and more energy-efficient deep-learning processors has therefore led to an intensive investigation of in-memory computation schemes using analog crossbar arrays.<sup>3,4</sup>

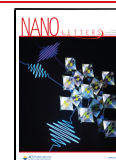
The basic building block of analog crossbar arrays is the crosspoint element, which can be described as a programmable, nonvolatile resistor.<sup>5</sup> For a DNN to be trained successfully (i.e., without degradation of classification performance) with such architectures, the resistive devices need to have many nonvolatile conductance states that can be modulated reversibly, symmetrically, and reproducibly.<sup>3,6</sup> A plethora of device technologies has been proposed for analog deep learning applications such as phase-change memories,<sup>7</sup> filamentary<sup>8</sup> and bulk-switching<sup>9</sup> resistive memories, ferroelectric tunnel junctions,<sup>10</sup> spintronics,<sup>11</sup> and superconducting nanowires.<sup>12</sup> Despite significant advances in the aforementioned fields, no device technology has been identified so far that meets all of the requirements.<sup>3</sup>

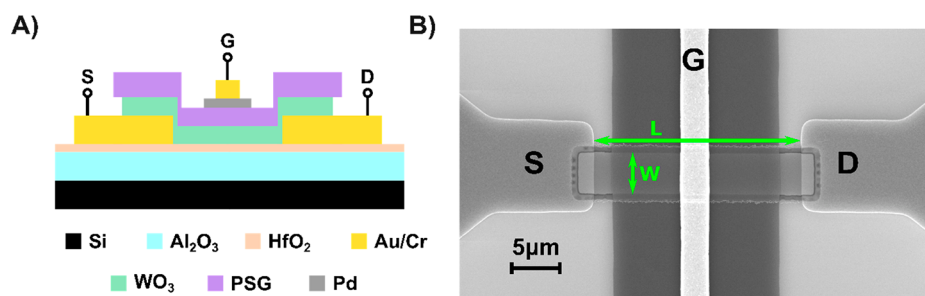
Recently, a novel device family relying on electrochemically controlled ion intercalation in transition-metal oxide channels (also referred to as electrochemical random access memory, ECRAM) has shown promising characteristics.<sup>13–19</sup> Due to their similarities to solid-state batteries, the most mature versions of these devices rely on shuffling Li<sup>+</sup> ions across a Li<sup>+</sup>-conducting electron-insulating electrolyte, metered by electron flow in the outer metallic circuit.<sup>13–15</sup> Approaches based on alkali-metal ions (Li<sup>+</sup>, Na<sup>+</sup>, K<sup>+</sup>, Ca<sup>2+</sup>) suffer from CMOS incompatibility, due to contamination concerns of the fabrication tools. Proton (H<sup>+</sup>) intercalation devices are thus considered to have the edge in realizing scalable programmable resistors that are back-end-of-line (BEOL) integrable.<sup>17–19</sup> Moreover, the smallest ionic radius of protons (Shannon–Prewitt effective ionic radius of  $-0.18 \text{ \AA}$  if 2-fold coordinated and  $-0.38 \text{ \AA}$  if one-fold coordinated) promises high modulation speed, low energy consumption, and enhanced endurance. Indeed, in our recent work,<sup>18</sup> we demonstrated promising device characteristics using a PdH<sub>x</sub> solid hydrogen

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**Figure 1.** (A) Schematic of the protonic programmable resistor studied in this work. (B) Top-view scanning electron microscope image of a finished device showing the source (S), drain (D), and gate (G) of a device with a nominal channel width ( $W$ ) of  $5\ \mu\text{m}$  and a length ( $L$ ) of  $25\ \mu\text{m}$ .

reservoir, polymeric Nafion electrolyte, and  $\text{WO}_3$  channel material for proton intercalation.

However, the main bottleneck with developing protonic programmable resistors has been the absence of a CMOS-compatible all-solid-state electrolyte that conducts protons but blocks electrons. All designs so far have relied on approaches that either cannot be integrated and scaled down, such as using organic materials,<sup>17</sup> using chemically and thermally sensitive polymers (e.g. Nafion),<sup>18</sup> or suffer from energy inefficiency such as high electric-field-induced water hydrolysis.<sup>19</sup> Furthermore, some protonic electrolytes rely on water absorption for proton conduction, requiring repeated rehydration of the material during operation.<sup>20</sup>

In this work, we demonstrate a CMOS-compatible protonic programmable resistor based on a nanoporous phosphosilicate glass (PSG) solid electrolyte. PSG is an excellent choice for this application due to its (1) good electronic insulation, (2) high room-temperature proton conductivity, (3) conduction mechanism based on P doping (instead of structural water), and (4) ready availability in conventional Si processing. The selection of this material has allowed us to use standard CMOS-fabrication techniques and consequently scale down the device footprint. Our devices show desirable modulation characteristics in terms of symmetry, retention, endurance, and energy efficiency. This new device technology appears promising in the quest to satisfy the stringent performance requirements for analog crosspoint elements. Furthermore, PSG enables exploration of alternative channel and hydrogen reservoir layers, while using a fully CMOS compatible and scalable fabrication platform.

## 2. DEVICE OPERATION

The protonic programmable resistor demonstrated in this work is a three-terminal device that employs a  $\text{WO}_3$  channel, a PSG electrolyte layer, and a Pd gate reservoir (Figure 1A). We adopt a transistor-like notation to refer to the three terminals of this device. The terms “source” and “drain” are used to refer to the two ends of the  $\text{WO}_3$  channel, while the term “gate” is used to refer to the  $\text{PdH}_x$  proton reservoir.

The basic operation principle of the device relies on modulating the channel conductance via the electrochemically controlled intercalation of protons into  $\text{WO}_3$ , as explained in our recent work.<sup>18</sup> Initially, protons are stored in the gate reservoir as  $\text{PdH}_x$ , which is achieved by the hydrogen uptake of Pd in a forming gas ambient (3%  $\text{H}_2$  in  $\text{N}_2$ ).<sup>21</sup> With the application of a voltage pulse, a controlled number of protons are moved from the gate to the channel through the solid electrolyte, while electrons concurrently move through the external circuit in the same direction. This process can be

reversed by the application of a voltage pulse of inverse polarity, forcing protons out of the channel back into the gate reservoir. Protons are n-type dopants in  $\text{WO}_3$ ,<sup>22</sup> and as they move in and out, the conductivity of the channel is modulated up and down. Furthermore, in the absence of a programming pulse (i.e., floating gate), proton motion toward either the gate or the channel is precluded, as the electrolyte prevents the electron flow required to satisfy charge neutrality in the channel and reservoir layers. This last feature establishes nonvolatility with the channel conductance remaining constant at its last programmed level.

Among several oxides whose electronic conductivity can be tuned via cation intercalation ( $\text{WO}_3$ ,<sup>18</sup>  $\text{V}_2\text{O}_5$ ,<sup>23</sup>  $\text{MoO}_3$ ,<sup>24</sup>  $\text{Nb}_2\text{O}_5$ ,<sup>25</sup>), we chose to use amorphous tungsten oxide (a- $\text{WO}_3$ ) as the channel material. This selection was motivated by the well-established conductivity modulation<sup>13,16,18</sup> and electrochromism<sup>19</sup> dynamics with cation intercalation. a- $\text{WO}_3$  is a CMOS-compatible semiconductor with a band gap of 2.8–3.2 eV whose conductivity can be precisely modulated by protonation, taking place concurrently with charge-balancing electron filling of the W 5d orbital dominated conduction band in the dilute regime. The structure of a- $\text{WO}_3$  at room temperature is assumed to be similar to that of its crystalline counterpart (monoclinic based on corner-sharing  $\text{WO}_6$  octahedra), but with disordered bond lengths and angles. The most common defect present in the  $\text{WO}_3$  lattice structure is an oxygen vacancy, which bonds to a  $\text{W}^{6+}$  ion, reduces the oxidation state of the neighboring  $\text{W}^{5+}$  ion, and increases the conductivity, in a way analogous to the way the electron does with proton intercalation (see Figure S1). The extent of conductivity modulation in a- $\text{WO}_3$  by proton intercalation, therefore, depends on its initial defect concentration that determines its initial conductivity.<sup>26</sup>

In the search for an inorganic CMOS-compatible electrolyte layer, we focused on silicate glasses ( $\text{SiO}_2$ ), which are widely used in Si technology as electron insulators.<sup>27</sup> Deposition conditions of silicate glasses can be engineered to yield a nanoporous structure with defect –OH terminated Si groups (silanol), providing a surface-site path for ion transport along the pores.<sup>28</sup> The acidic nature of silanol acts as a proton donor, which then migrates by hopping between hydroxyl groups and structural water.<sup>29,30</sup> Doping silicate glasses with phosphorus sterically hinders the glass network to increase nonbridging oxygen bonds,<sup>31</sup> replaces Si–O–Si bonds with –Si–OH and –Si–O–P–OH groups, and increases both the pore volume and surface area.<sup>32</sup> The P–OH groups not only have higher acidity in comparison to silanol but are also amphoteric, meaning that they can act as both proton donors and acceptors.<sup>33</sup> These are all key properties that provide

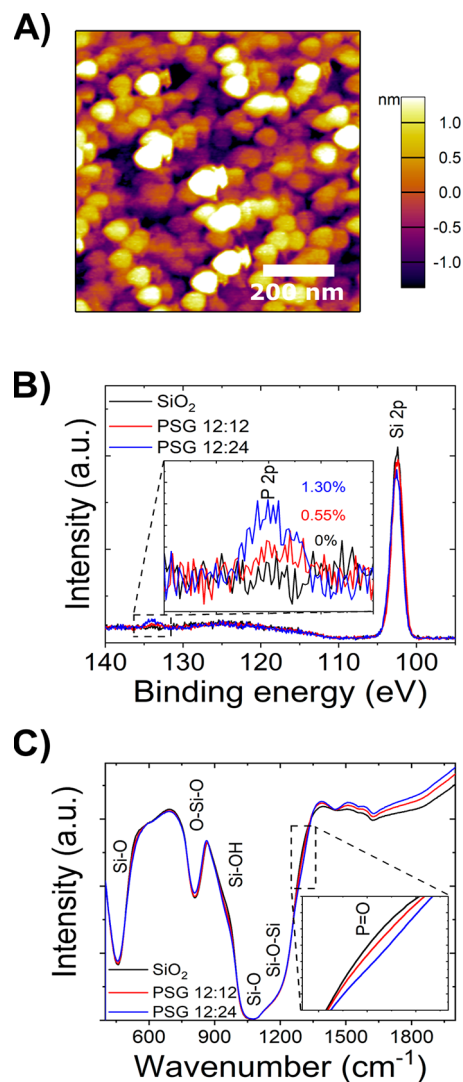
phosphosilicate glass (PSG, P-SiO<sub>2</sub>) high proton conductivity at room temperature ( $2.54 \times 10^{-4}$  S cm<sup>-1</sup>), which stands among the highest values in comparison with several perovskite, fluorite, and simple oxide proton-conducting materials,<sup>34</sup> while retaining their electron-insulating properties.<sup>35</sup> Furthermore, even though the proton conductivity of PSG can be further increased to 10<sup>-1</sup> S cm<sup>-1</sup> by an additional hydrothermal treatment<sup>36</sup> and chemisorption of water in the pore structure,<sup>37</sup> such modifications also result in an increased electronic conductivity of the material, which is undesirable for our design that requires nonvolatile throttling of hydrogen in WO<sub>3</sub>.

### 3. METHODS

The protonic programmable resistors were fabricated on a Si substrate covered with 10/90 nm HfO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> deposited by atomic layer deposition (ALD) for electrical and protonic insulation. Channel contacts (15/5 nm Au/Cr) were first patterned using a direct-write photolithography and liftoff process. A 10 nm WO<sub>3</sub> channel and 10 nm PSG electrolyte layers were blanket-deposited using ALD and plasma-enhanced chemical vapor deposition (PECVD) processes, respectively. The deposition conditions for PSG ( $T = 100$  °C, 60 W plasma power, and gas flow ratio of 12 sccm SiH<sub>4</sub>/12 sccm PH<sub>3</sub> diluted at 2% in H<sub>2</sub>) were found to be optimal to maximize proton conductivity, as reported in ref 35. The PSG/WO<sub>3</sub> stack was subsequently patterned with a self-aligned reactive ion etching (RIE) process, followed by TMAH-based wet etching of WO<sub>3</sub> to create a PSG overhang that prevents shorting of the channel with the gate at the edges of the device. Different channel dimensions (i.e., width and length) were patterned in the range between 2 and 100 μm. Finally, the 5 nm Pd reservoir and 150/10 nm Au/Cr gate interconnect and pads were electron beam evaporated and patterned through separate liftoff processes. A cross-sectional view of a finished device is shown in Figure 1A. Full details of the fabrication process can be found in Figure S2 in the Supporting Information. Figure 1B shows the top-view scanning electron microscope image of a fabricated structure. Given the standard Si-fabrication compatibility of this process flow, it can easily be adapted to yield nanoscale devices by replacing the photolithography steps with their electron-beam lithography versions.

An atomic force microscopy (AFM) image of a PSG film deposited under the optimum conditions on a Si substrate is shown in Figure 2A. A nanogranular structure (so-called nanoglass) with a mean glassy grain diameter of ~80 nm and an RMS roughness of ~1 nm is observed. This image also evidences the presence of nanopores with a high surface to volume ratio, an essential requirement for efficient proton transport in this material.

The stoichiometric ratio of each element in the optimized PSG film (PSG 12:12) was found to be 0.55% P, 45.65% Si, and 53.8% O using X-ray photoelectron spectroscopy (XPS). The spectra of undoped (SiO<sub>2</sub>), optimized (PSG 12:12), and more highly doped (PSG 12:24) films, deposited with 12:0, 12:12 and 12:24 SiH<sub>4</sub>:PH<sub>3</sub> flux, respectively, were acquired (Figure 2B). While the intensity of the P 2p peak, located at a binding energy of ~134 eV (inset), increases with PH<sub>3</sub> flux, that of the Si 2p peak decreases. The estimated P concentration of the PSG 12:12 film is 0.55%, a value similar to those reported for PSG films with optimized proton conductivity ( $2.5 \times 10^{-4}$  S cm<sup>-1</sup>).<sup>35</sup>



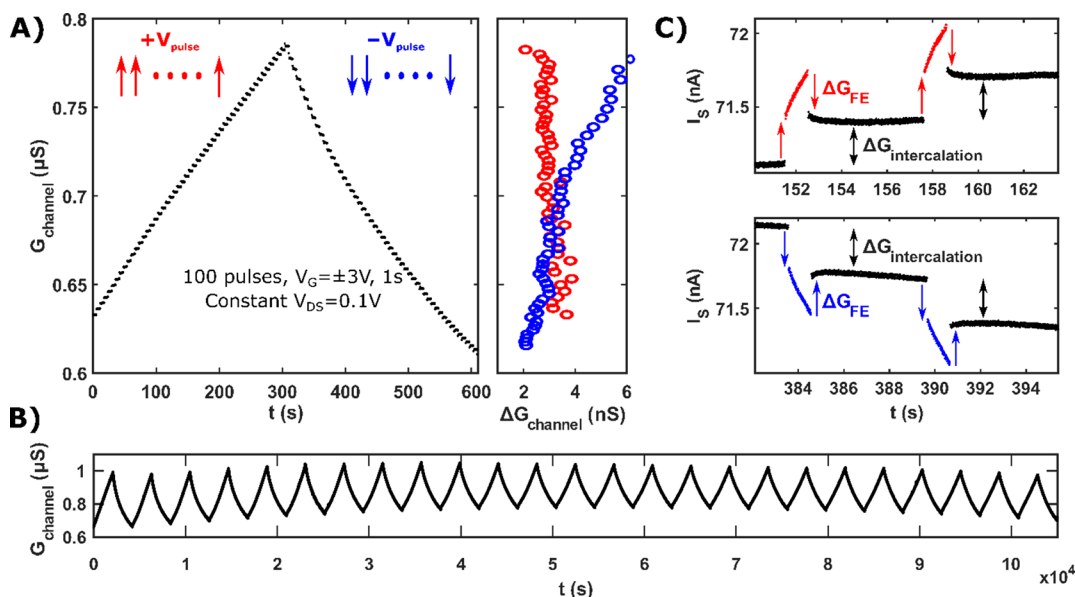
**Figure 2.** (A) Atomic force microscopy image of the PSG 12:12 thin film surface deposited on a Si surface. (B) X-ray photoelectron spectra of the P 2p and Si 2p peaks and (C) Fourier transform infrared spectra of the SiO<sub>2</sub>, PSG 12:12, and PSG 12:24 thin films. The inset in (B) shows enlarged spectra in the P 2p energy range, while the inset in (C) shows enlarged spectra in the P=O absorption region.

The presence of P in the films was also confirmed by Fourier transform infrared spectroscopy (FTIR). As shown in Figure 2C, the increase in PH<sub>3</sub> flux results in the appearance of P=O and Si-OH absorption bands around 1300 cm<sup>-1</sup> (inset) and 1130 cm<sup>-1</sup>, respectively, as well as a decrease in the intensity of the Si-O and Si-O-Si bands. All of these features are indicative of the successful incorporation of P atoms in the glass structure.

### 4. RESULTS

Electrical characterization of the devices was conducted at room temperature in an enclosed probe station (NEXTRON MPS-PT) filled with forming gas (3% H<sub>2</sub> in N<sub>2</sub>, see the Supporting Information for further details). The basic modulation properties of the devices were first measured to assess their conductance modulation depth, symmetry, retention, and endurance (Figure 3) in a pulsing mode. Figure 3A shows the increase in conductance upon the application of 100 positive voltage pulses (increment, red,  $V_G = 3$  V, 1 s) and





**Figure 3.** (A) (left) Modulation characteristics of a protonic programmable resistor ( $5 \mu\text{m}$  width,  $50 \mu\text{m}$  length) determined by applying 50 voltage pulses of  $\pm 3 \text{ V}$  and  $1 \text{ s}$  in either direction of the gate voltage. Each state is monitored for  $2 \text{ s}$  by measuring the channel current while  $V_{\text{DS}} = 0.1 \text{ V}$  is applied. (right) Conductance change ( $\Delta G_{\text{channel}}$ ) plotted at different conductance levels ( $G_{\text{channel}}$ ). Incremental changes are shown with red markers, while decremental changes are shown with blue markers. (B) Endurance characterization throughout the application of 50000 voltage pulses ( $25 \times [1000 \uparrow 1000 \downarrow]$ ) of  $\pm 3 \text{ V}$  and  $0.1 \text{ s}$  width over the course of  $\sim 30 \text{ h}$ . (C) Detailed picture of (A) where the channel current,  $I_{\text{S}}$ , is constantly recorded in the presence (colored dots) and absence (dark dots) of gate pulses.

decrease upon that of 100 negative pulses (decrement, blue,  $V_{\text{G}} = 3 \text{ V}$ ,  $1 \text{ s}$ ).  $V_{\text{DS}}$  was constantly maintained at  $0.1 \text{ V}$  throughout the update sequence.

Figure 3A demonstrates very clean device conductance modulation characteristics with high symmetry and non-volatility of the programmed conductance, due to the electron-blocking properties of the PSG thin-film electrolyte (see Figure S3). A well-defined symmetry point (i.e., a conductance level at which the incremental change is equal to the decremental change in conductance) is also evidenced, as well as excellent cycling endurance characteristics, with very little variation in conductance values even after the application of 50000 pulses over the course of  $\sim 30 \text{ h}$  (Figure 3B). Furthermore, we did not observe any noticeable open-circuit voltage development across the gate stack, allowing the operation under constant voltage pulses.

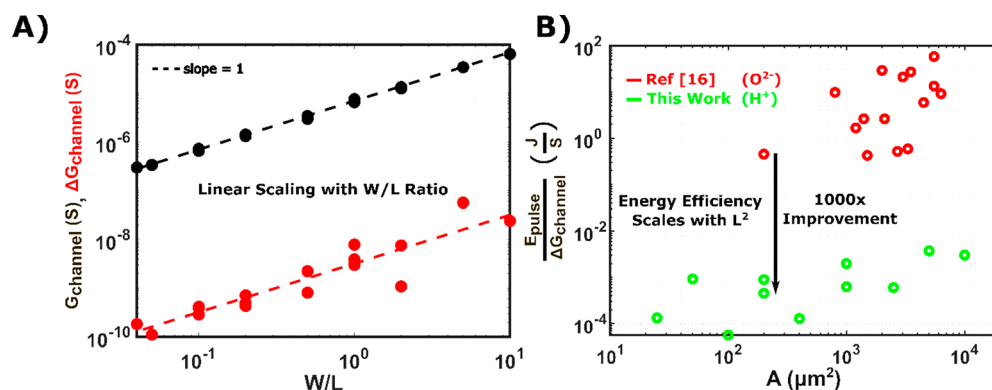
It should be noted that the base (i.e., unprotonated) conductance of the channel material is significantly higher than that of stoichiometric  $\text{WO}_3$ , therefore limiting the conductance modulation depth and the speed of the channel upon proton intercalation.<sup>18,26</sup> Such behavior is expected on consideration of the oxygen-deficient nature of the ALD-deposited  $\text{WO}_3$ . Indeed, an XPS analysis (Figure S6) clearly evidences a large degree of initial tungsten reduction ( $\text{W}^{5+}:\text{W}^{6+}$  ratio) and in-gap states present in the film. To mitigate this issue, the deposition method of  $\text{WO}_3$  can be modified to target the optimal stoichiometry or alternative channel materials (e.g.,  $\text{V}_2\text{O}_5$ ,  $\text{MoO}_3$ ,  $\text{Nb}_2\text{O}_5$ ) can be used as a replacement.

A careful analysis of the modulation dynamics (Figure 3C) unravels interesting device physics, comprised of both volatile and nonvolatile changes in conductance, as programming pulses are applied. Under constant  $V_{\text{DS}} = 0.1 \text{ V}$ , the source current,  $I_{\text{S}}$ , immediately steps up/down when a gate pulse ( $\pm 3 \text{ V}$ ) is applied and then smoothly increases/decreases for the remainder of the pulse duration (colored dots). Once the pulse

disappears (i.e., floating gate), another sudden drop/rise in  $I_{\text{S}}$  is observed. In the absence of a gate pulse,  $I_{\text{S}}$  remains approximately constant, at a level different from that before the application of the pulse, indicating nonvolatile programming of the channel conductance (dark dots). These sudden changes in  $I_{\text{S}}$  do not reflect an increase in gate current,  $I_{\text{G}}$  ( $< \text{pA}$ ), but are due to a field-effect enhancement of the channel conductance ( $\Delta G_{\text{FE}}$ ). This behavior can be explained by a field-effect increase in the electron concentration and a resulting increase in conductance of the n-type  $\text{WO}_3$  channel by the electrostatic field of protons driven within the electrolyte close to the  $\text{WO}_3$  interface.<sup>38,39</sup> Unlike the nonvolatile, electrochemical intercalation induced conductance modulation ( $\Delta G_{\text{intercalation}}$ ), this additional channel current only flows during the application of a gate voltage pulse and therefore it is volatile.

The initial channel conductance and average modulation magnitude of protonic crosspoint elements for 20 devices with varying dimensions were also examined. Figure 4A shows that both parameters scale linearly with the channel width/length ( $W/L$ ) ratio, as expected. This well-behaved scaling demonstrates both a promising yield and uniformity for these PSG-based devices.

Using different channel dimensions, we have also estimated the energy efficiency and its scaling for protonic programmable resistors presented in this work. As the figure of merit, we chose the ratio of energy consumption required to change the channel conductance by a small amount ( $E/\Delta G_{\text{channel}}$ ). This choice also enables us to compare the same value with those reported for  $\text{O}^{2-}$  ion intercalation devices in ref 16. Figure 4B reveals that, at the same device area of  $200 \mu\text{m}^2$ , our devices consume  $0.45 \text{ mJ/S}$ , which is  $1000\times$  more energy efficient with respect to its heavier-ion alternative. Considering that this metric scales with  $L^2$  (not area), we estimate the pulse energy required for a single state increment of a  $100 \times 100 \text{ nm}^2$



**Figure 4.** (A) Scaling of conductance and average conductance modulation magnitude for different channel geometries ( $W/L$ ). (B) Comparison of device energy efficiency with respect to other CMOS-compatible oxygen ion intercalation devices reported in ref 16.

protonic device would be  $\sim 1$  fJ, assuming a conductance range of 1–10  $\text{M}\Omega$  modulated with 1 nS incremental changes (i.e., 900 states). We expect this number to further improve upon the implementation of a stoichiometric  $\text{WO}_3$  channel, since the oxygen-deficient nature of the current material limits its overall modulation amplitude ( $\Delta G_{\text{channel}}$ ), as explained earlier.

## 5. CONCLUSION

In this work, we demonstrated a back-end CMOS-compatible protonic programmable resistor that can serve as a crosspoint element in analog crossbar processors. Conductance modulation, device scalability, and process control are ensured by the use of a thin nanoporous PSG layer, a common and CMOS-compatible material, as the proton electrolyte layer. Our devices, which also include a  $\text{WO}_3$  channel and a  $\text{PdH}_x$  reservoir, display low energy consumption along with symmetric and nonvolatile conductance modulation for an extended number of voltage pulses. Downscaling of the device dimensions, using electron-beam lithography or other techniques, as well as the possibility of BEOL integration to CMOS chips, can be also envisioned for our technology. Ultimately, PSG can serve as a platform to explore alternative channel and hydrogen reservoir layers in protonic programmable resistors, which present desirable characteristics to realize next-generation analog accelerators for deep-learning applications.

## ■ ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.nanolett.1c01614>.

Details of fabrication, characterization methods, additional results, and discussions (PDF)

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## Notes

The authors declare no competing financial interest.

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# Supplementary Materials for CMOS-compatible Protonic Programmable Resistor based on Phosphosilicate Glass Electrolyte for Analog Deep Learning

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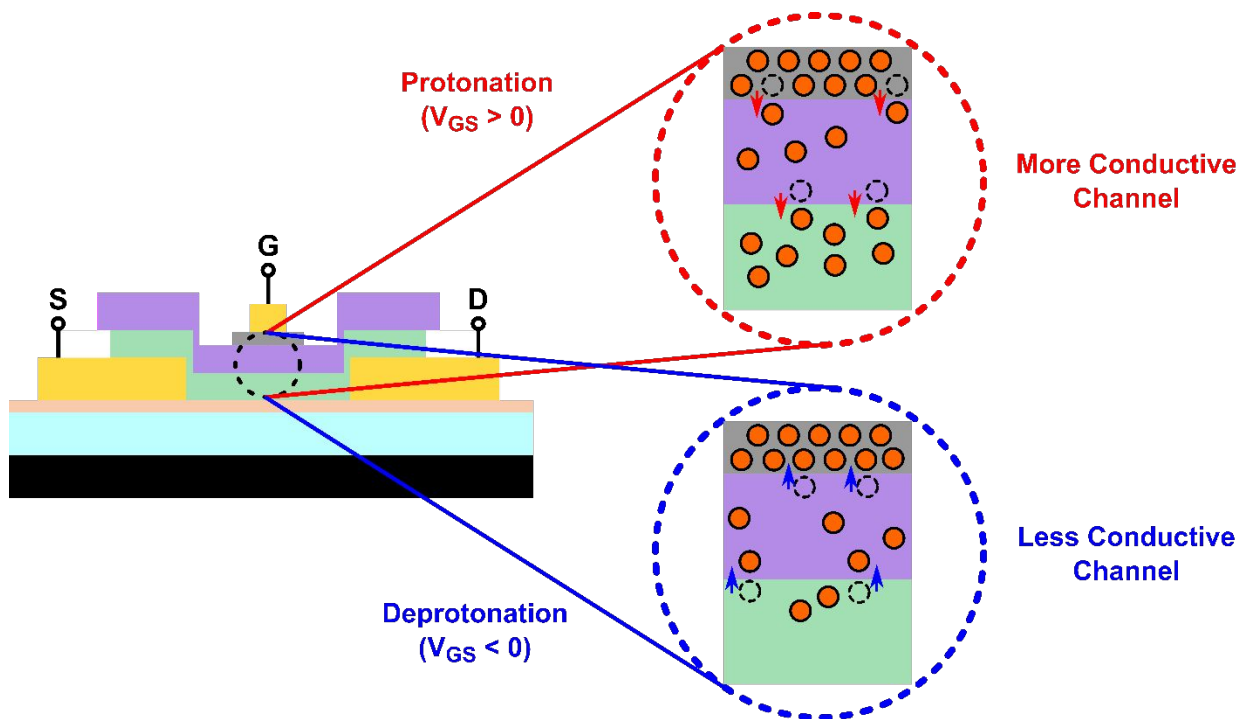
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## 1. Device Operation



**Figure S1.** Operational illustration of CMOS-compatible protonic nonvolatile resistors studied in this work.

## 2. Fabrication Flow

- Atomic Layer Deposition (ALD) of 90 nm  $\text{Al}_2\text{O}_3$  (for proton blocking properties) and 10 nm  $\text{HfO}_2$  (to protect  $\text{Al}_2\text{O}_3$  from TMAH and  $\text{CF}_4$  plasma) on an undoped 4" Si wafer for electronic and protonic insulation.
- Spin-coating of Polymethylglutarimide (PMGI) at 4 krpm for 60 s (thickness  $\approx 120$  nm) and baking at  $180^\circ\text{C}$  for 90 s. Spin-coating of Microposit S1813 positive tone photoresist at 4 krpm for 60 s (thickness  $\approx 1.8 \mu\text{m}$ ) and baking at  $100^\circ\text{C}$  for 90 s.
- Exposition of the source and drain electrode (M1) layout using a Heidelberg maskless aligner (MLA) 150 with a 405 nm laser wavelength and at a dose of  $80 \text{ mJ/cm}^2$ .
- Development of the photoresist in MF CD-26 for 80 s followed by rinsing in deionized (DI) water for 60 s.
- Deposition of 5 nm Cr and 15 nm Au using an electron-beam evaporator.
- Liftoff by soaking the wafer in N-methyl pyrrolidone (NMP) for 5 hours.
- Removal of excess PMGI by soaking the wafer in MF CD-26 for 30 s.
- Spin-coating of Microposit S1813 positive tone photoresist at 4 krpm for 60 s (thickness  $\approx 1.8 \mu\text{m}$ ) and baking at  $100^\circ\text{C}$  for 90 s.
- Dicing of the 4" wafer into  $1 \times 1 \text{ cm}^2$  samples.
- Cleaning of the samples by sonication in acetone and isopropanol baths for 300 s each.
- Deposition of 10 nm  $\text{WO}_3$  using ALD with Bis(tert-butylimino)bis(dimethylamino)tungsten (VI) (BTBMW) and  $\text{O}_3$  precursors at  $330^\circ\text{C}$ .
- Plasma-Enhanced Chemical Vapor Deposition (PECVD) of 10 nm PSG using 1420 sccm  $\text{N}_2\text{O}$ , 12 sccm  $\text{SiH}_4$ , and 12 sccm  $\text{PH}_3$  (2% in  $\text{H}_2$ ) at  $100^\circ\text{C}$ , with a RF plasma power of 60 W at 380 kHz.
- Spin-coating of Microposit S1813 positive tone photoresist at 4 krpm for 60 s (thickness  $\approx 1.8 \mu\text{m}$ ) and baking at  $100^\circ\text{C}$  for 90 s.
- Exposition of the active layer ( $\text{WO}_3$  and PSG channel pattern) layout using a Heidelberg-MLA 150 with a 405 nm laser wavelength and at a dose of  $80 \text{ mJ/cm}^2$ .
- Development of the photoresist in MF CD-26 for 60 s followed by rinsing in DI water for another 60 s.
- Patterning of both PSG and  $\text{WO}_3$  layers using Reactive Ion Etching (RIE) with a  $\text{CF}_4$  plasma at 100 W for 180 s (partitioned in 60 s intervals).
- Selective wet-etching (undercut) of the  $\text{WO}_3$  layer in MF CD-26 for 300 s at room temperature.
- Stripping of the excess photoresist in N-methyl pyrrolidone (NMP) bath for 12 hours.



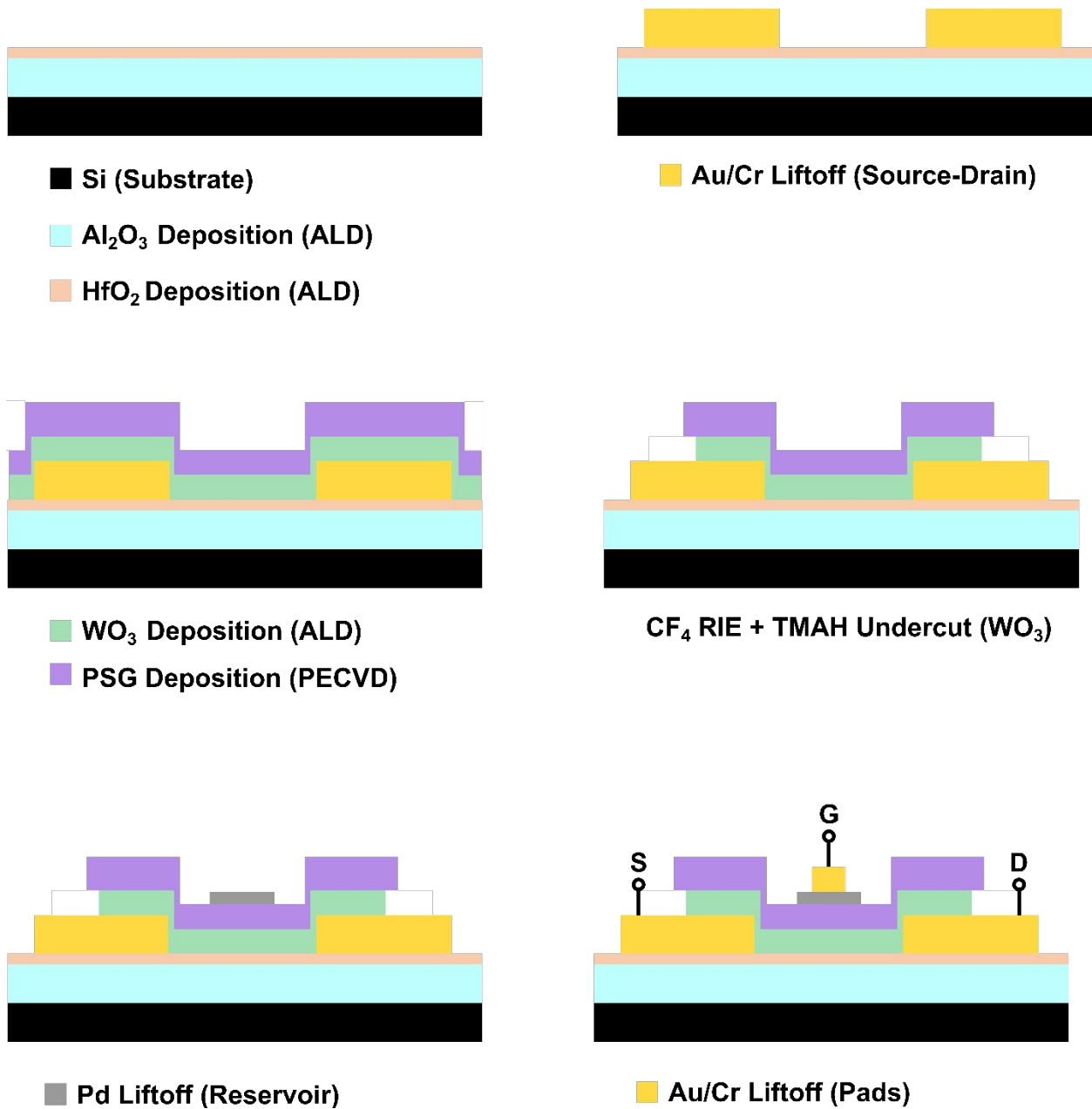
- Spin-coating of PMGI at 4 krpm for 60 s (thickness  $\approx 120$  nm) and baking at  $180^\circ\text{C}$  for 90 s. Spin-coating of Microposit S1813 positive tone photoresist at 4 krpm for 60 s (thickness  $\approx 1.8\ \mu\text{m}$ ) and baking at  $100^\circ\text{C}$  for 90 s.
- Exposition of the gate electrode layout (M2) using Heidelberg-MLA 150 with a 405 nm laser wavelength and at a dose of  $80\ \text{mJ}/\text{cm}^2$ .
- Development of the photoresist in MF CD-26 for 80 s followed by rinsing in DI water for another 60 s.
- Deposition of 5 nm Pd using an electron-beam evaporator at a low deposition rate ( $\leq 0.5\ \text{\AA}/\text{s}$ ).
- Liftoff by soaking the pieces in NMP for 5 hours.
- Removal of the excess PMGI by soaking the pieces in MF CD-26 for 30 s
- Spin-coating of PMGI at 4 krpm for 60 s (thickness  $\approx 120$  nm) and baking at  $180^\circ\text{C}$  for 90 s. Spin-coating of Microposit S1813 positive tone photoresist at 4 krpm for 60 s (thickness  $\approx 1.8\ \mu\text{m}$ ) and baking at  $100^\circ\text{C}$  for 90 s.
- Exposition of the pad (M3) layout using Heidelberg-MLA 150 with a 405 nm laser wavelength and at a dose of  $80\ \text{mJ}/\text{cm}^2$ .
- Development of the photoresist in MF CD-26 for 80 s followed by rinsing in DI water for 60 s.
- Deposition of 10 nm Cr and 150 nm Au using an electron-beam evaporator.
- Liftoff by soaking the wafer in NMP for 5 hours.

The sheet resistance of the ALD  $\text{WO}_3$  was  $\sim 18.2\ \text{k}\Omega/\square$  as deposited, which then slightly decreased to  $\sim 13.5\ \text{k}\Omega/\square$  during the fabrication process. On the other hand, the PSG layer we used in this work shows good electronic insulation properties. For example, for a device with 5  $\mu\text{m}$  width and 50  $\mu\text{m}$  length, 3 V gate voltage yielded a current of  $\sim 200\ \text{pA}$  ( $\sim 15\ \text{G}\Omega$ ), indicating that the electrolyte layer is much more resistive than the channel layer ( $R_{\text{SD}} \sim 1.4\ \text{M}\Omega$ ).

Regarding the yield rate of this process, 130 devices were fabricated on a  $1 \times 1\ \text{cm}^2$  chip. Out of the 50 devices we measured, 41 functioned successfully and reproducibly over the course of more than 2 months of experimentation.

Since we demonstrate standard Si fabrication compatibility, the fabrication flow can be easily adapted to yield nanoscale devices by simply replacing photolithography steps with electron-

beam lithography steps.



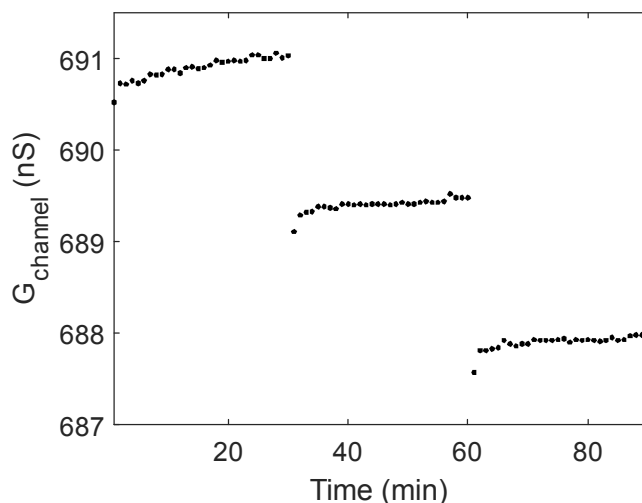
**Figure S2.** Fabrication flow of CMOS-compatible protonic nonvolatile resistors based on PSG electrolyte.

### 3. Methods

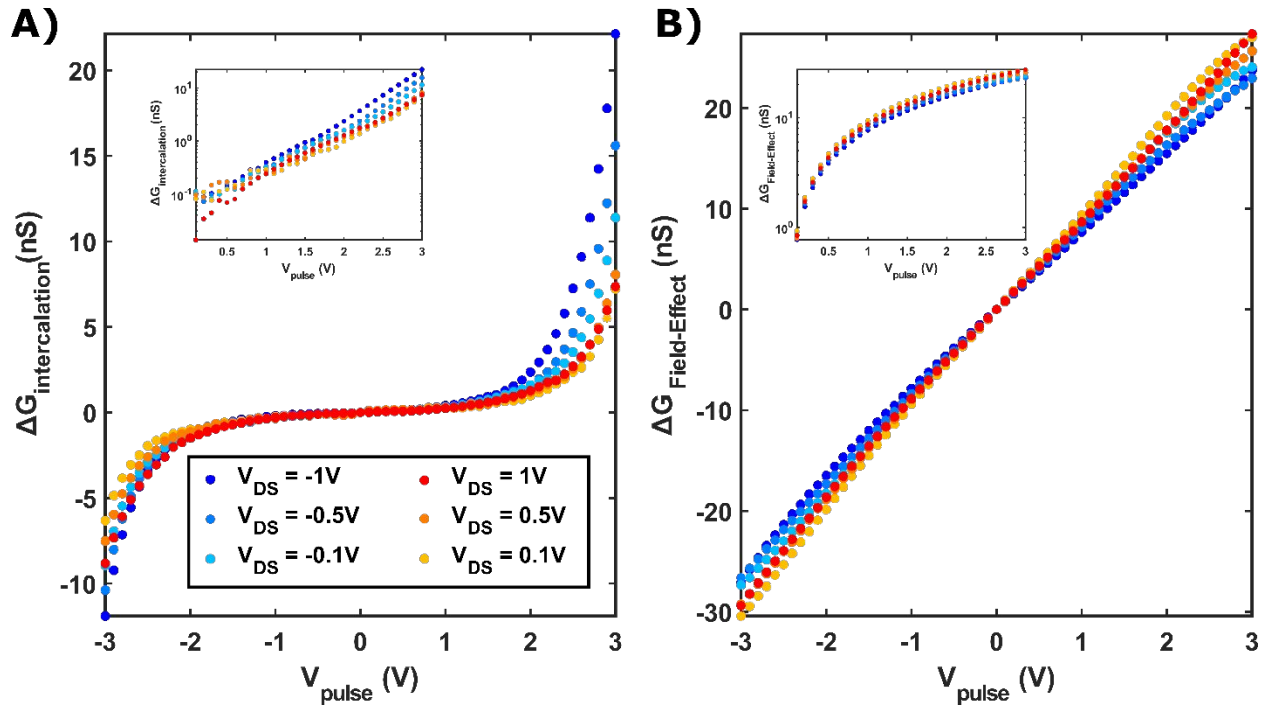
All electrical measurements were performed using a micro-probe station enclosed chamber (MPS-PT) manufactured by NEXTRON, Korea. In experiments performed in forming gas (FG, 3% H<sub>2</sub> in N<sub>2</sub>) all 4 DC probes were first connected to the pads of the device under test (1 to Source, 1 to Drain and 2 to Gate terminals). The chamber was then purged by flowing in FG for 60 s with both gas inlet and outlet open, followed by another 60 s with outlet shut to create positive pressure inside the chamber.

Three of the probes (1 Source, 1 Drain and 1<sup>st</sup> Gate) were then connected to the Source Measurement Units (SMUs) of a Keysight B1500 Semiconductor Analyzer, while the fourth probe (2<sup>nd</sup> Gate) was connected to the Pulse Generation Unit (PGU) of the same instrument. The experiment sequence and data acquisition were controlled via an in-house developed MATLAB suite.

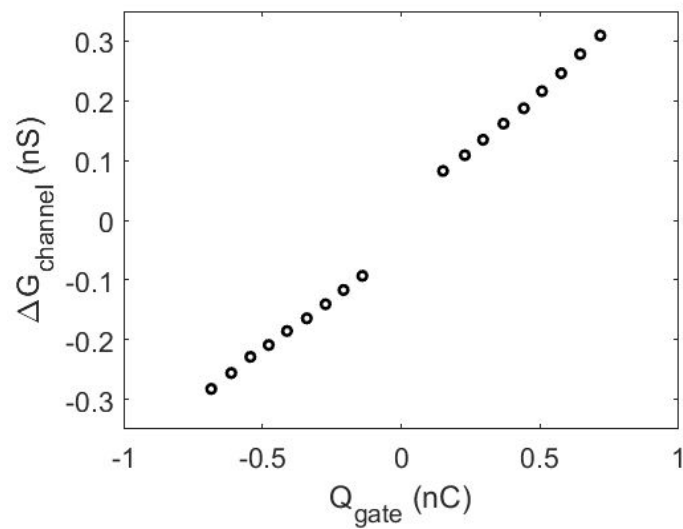
### 4. Additional Results



**Figure S3.** Retention characteristics of the protonic device incrementally programmed by a voltage pulse ( $V_{pulse} = -4 V$ ,  $t_{pulse} = 1 s$ ) at 30 and 60 minutes. Each conductance state is read every minute by sweeping  $V_{SD}$  between  $\pm 0.1 V$  while the gate terminal is floating. Conductance of the device remains distinctly separate from adjacent levels over 30 minutes of readout.

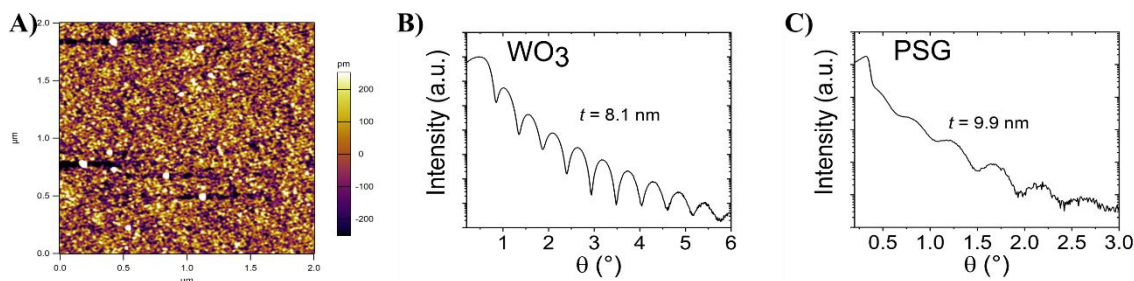


**Figure S4.** (A) Intercalation and (B) field-effect conductance modulation as a function of the pulse gate voltage for different read drain-source voltage.



**Figure S5.** Linear dependence of change in channel conductance ( $\Delta G_{\text{channel}}$ ) as a function of charge injected to / extracted from the gate ( $\Delta Q_{\text{gate}}$ ).

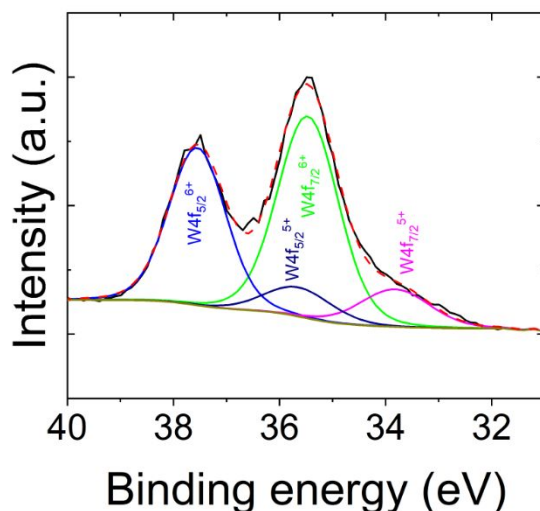




**Figure S6.** (A) Atomic force microscopy image of the reference ALD  $\text{WO}_3$  thin film surface, deposited on a Si surface alongside the patterned sample. X-ray reflectivity (XRR) curve of the reference (B) ALD  $\text{WO}_3$  and (C) PSG 12:12 thin films, deposited on a Si surface alongside the patterned samples.

The very smooth surface and low RSM roughness ( $\approx 0.2$  nm) of the ALD  $\text{WO}_3$  film as well as the presence of well-defined thickness-interference (Kiessig) fringes in the XRR curves of both  $\text{WO}_3$  and PSG indicates that the films' surface is smooth and the film-substrate interface well defined. The film thickness ( $t$ ) can thus be calculated from these XRR curves from the following equation  $t = \frac{(m-n)\lambda}{2(\sin\theta_m - \sin\theta_n)}$

where  $m$  and  $n$  are the orders of interference, and  $\theta_m$  and  $\theta_n$  are the corresponding diffraction angles, respectively.



**Figure S7.** X-ray photoelectron spectra of the W  $4f$  peak of a  $\text{WO}_3$  thin film deposited by ALD. The W  $4f_{5/2}$  and W  $4f_{7/2}$  peaks deconvolution between the  $\text{W}^{6+}$  and  $\text{W}^{5+}$  oxidation states evidences a strong  $\text{W}^{5+}$  contribution ( $\approx 1 \text{ W}^{5+} : 4 \text{ W}^{6+}$  ratio), indicative of the initial reduced (substoichiometric,  $\text{WO}_{3-x}$ ) nature of the film.

## 5. Additional Discussion

Unlike the P-concentration in PSG, we did not control the amount of H in Pd. From the pressure-concentration isotherms of the H/Pd system given by Lewis (Platinum Metals Rev., **4**, 132, 1960), the atomic ratio of H/Pd at room temperature and ambient pressure is  $\sim 0.65$ . Since we use premixed 3% H<sub>2</sub> in N<sub>2</sub>, we have only measured devices with this H/Pd ratio and assumed that the amount of H shuffled in/out of the WO<sub>3</sub> channel is much less than the amount of H stored in the reservoir. There are two reasons behind running the experiments under forming gas environment. (1) Under storage conditions (N<sub>2</sub>) prior to experimentation, Pd has no H within. The hydrogenation occurs when the device is enclosed in a forming gas ambient. (2) For Pd to retain H, it should not be exposed to O<sub>2</sub>, which causes hydrogen to oxidize and leave the reservoir material. Once the reservoir is loaded with H, it can alternatively be operated under an inert environment (e.g. N<sub>2</sub>, Ar, vacuum) as well.

We want to note that in measurements performed on devices with thicker Pd layers ( $\sim 80$  nm), we observed that the H uptake caused volume expansion of the Pd layer to the extent of it getting exfoliated. We anticipate that higher H<sub>2</sub> concentration could therefore be problematic even for devices with thinner Pd layers. On the other hand, as the number of protons that are exchanged with the channel layer is much less than what is stored in the reservoir, we expect that devices would be operational under lower H<sub>2</sub> concentrations as well.